

# Third-order accurate MPDATA for variable flows

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# Outline of the talk

1) Motivation

2) Main ideas of the scheme derivation

3) Benefits of the fully third-order MPDATA for global tracer transport

4) Benefits of the fully third-order MPDATA for fluid dynamics applications

5) Conclusions and outlook

## General motivation for high-order schemes

- Using a high-order scheme can be more cost-effective than increasing resolution with a lower order scheme.
- The functional form of a scheme leading-order truncation error determines its behavioural aspects such as respecting solution symmetries and minimising implicit diffusion or dispersion.
- High-order schemes with compact stencils make more efficient use of modern computers due to usually higher arithmetic intensity than lower order methods.

## Motivation for third-order MPDATA for variable flows

- Third-order MPDATA for constant flows was derived in L. Margolin, P. K. Smolarkiewicz, SISC, 1998.
- This variant exhibits more uniform distribution of the truncation error as function of the Courant number leading to better preservation of the solution symmetries.
- It is especially beneficial for tracer transport.

## Background

- MPDATA is a sign-preserving second-order accurate scheme for numerical integration of the generalised transport equation

$$\frac{\partial G\Psi}{\partial t} + \nabla \cdot (\mathbf{V}\Psi) = 0$$

based on iterative application of the first-order-accurate upwind algorithm.

- The basic version with two iterations uses physical velocity  $\mathbf{V}$  in the first iteration and leading-error-compensating pseudo-velocity  $\overline{\mathbf{V}}$  in the second.
- Pseudo-velocity is derived via truncation error analysis of the upwind scheme.

## Pseudo-velocity on structured grids

$$\overline{\mathbf{V}} = \frac{1}{2} \delta \mathbf{x} \odot \uparrow \mathbf{V} \uparrow \odot \frac{\nabla \Psi}{\Psi} - \frac{1}{2} \delta t \frac{\mathbf{V}}{G} \left[ \frac{\nabla \cdot (\mathbf{V}\Psi)}{\Psi} \right]$$

where

$$(\uparrow \mathbf{a} \uparrow)^I = |a^I|$$

denotes component-wise absolute value of a vector and

$$(\mathbf{a} \odot \mathbf{b})^I := a^I b^I$$

is the Hadamard product of two vectors

# Third-order accurate MPDATA for time and space dependant flows

## New idea

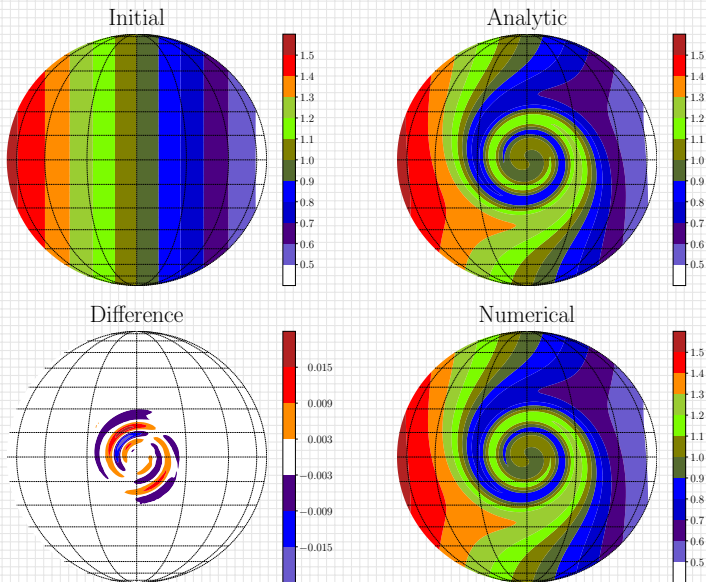
- Find a new pseudo-velocity  $\overline{\overline{\mathbf{V}}}$  that compensates the leading-order MPDATA error for variable flows.
- The derivation of  $\overline{\overline{\mathbf{V}}}$  uses MPDATA as a starting point in contrast to the derivation of the constant coefficients third-order scheme.
- Analytical calculation were verified and extended by using symbolic computer algebra.

## Third-order error-compensating velocity

$$\begin{aligned}\overline{\overline{\mathbf{V}}} = & -\frac{\delta \mathbf{x} \odot \delta \mathbf{x}}{24} \odot \left[ 4\mathbf{V} \odot \frac{1}{\Psi} \nabla \odot \nabla \Psi + 2\frac{\nabla \Psi}{\Psi} \odot \nabla \odot \mathbf{V} + \alpha \nabla \odot \nabla \odot \mathbf{V} \right] \\ & + \beta_M \frac{\delta \mathbf{x}}{2} \odot \uparrow \overline{\mathbf{V}} \uparrow \odot \frac{\nabla \Psi}{\Psi} + \frac{\delta t}{2} \delta \mathbf{x} \odot \uparrow \mathbf{V} \uparrow \odot \frac{1}{\Psi} \nabla \left[ \frac{1}{G} \nabla \cdot (\mathbf{V} \Psi) \right] \\ & + \frac{\delta t^2}{24} \left\{ -\frac{8\mathbf{V}}{G\Psi} \nabla \cdot \left[ \frac{\mathbf{V}}{G} \nabla \cdot (\mathbf{V} \Psi) \right] + \gamma \frac{\partial^2 \mathbf{V}}{\partial t^2} + \frac{2\mathbf{V}}{G\Psi} \nabla \cdot \left( \frac{\partial \mathbf{V}}{\partial t} \Psi \right) - \frac{2}{G\Psi} \frac{\partial \mathbf{V}}{\partial t} \nabla \cdot (\mathbf{V} \Psi) \right\}\end{aligned}$$

# Benefits of the fully third-order MPDATA for global tracer transport

# Moving vortices on the sphere – Nair & Jablonowski MWR 2008



## Sign-preserving MPDATA variants

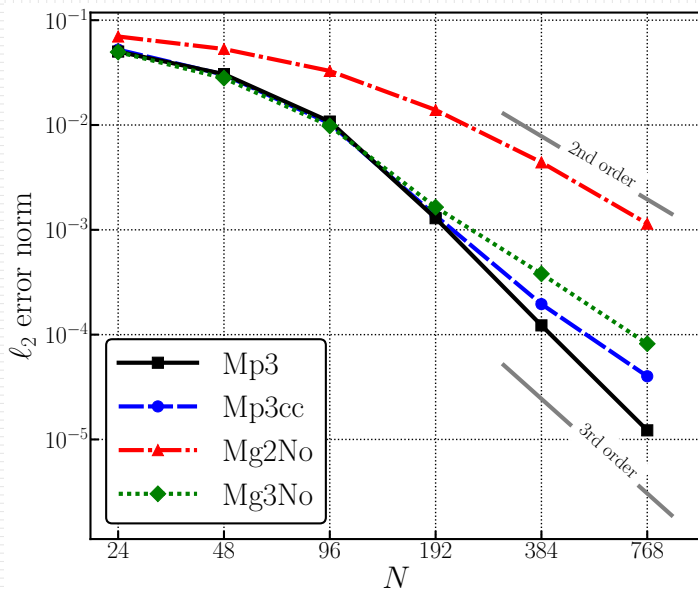
- Mp2 – fully second-order-accurate MPDATA
- Mp3 – fully third-order-accurate MPDATA
- Mp3cc – third-order constant-coefficient MPDATA

## Nonoscillatory infinite-gauge MPDATA variants

- Mg2No – nonoscillatory infinite-gauge variant of Mp2
- Mg3No – nonoscillatory infinite-gauge variant of Mp3
- Mg3ccNo – nonoscillatory infinite-gauge variant of Mp3cc



# Moving vortices on the sphere – convergence results



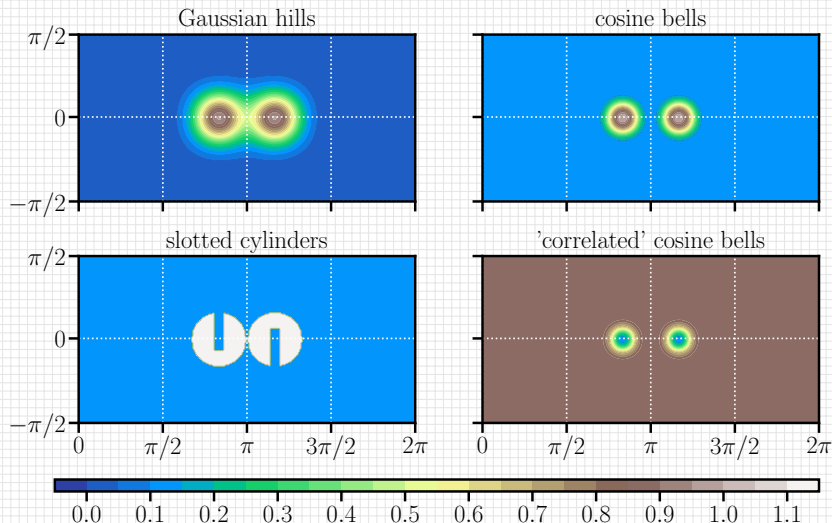
## Selected runtimes relative to the upwind scheme

Upwind	Mp2	Mg2No	Mp3cc	Mp3	Mg3No
1.0	3.6	5.9	9.5	10.3	12.6

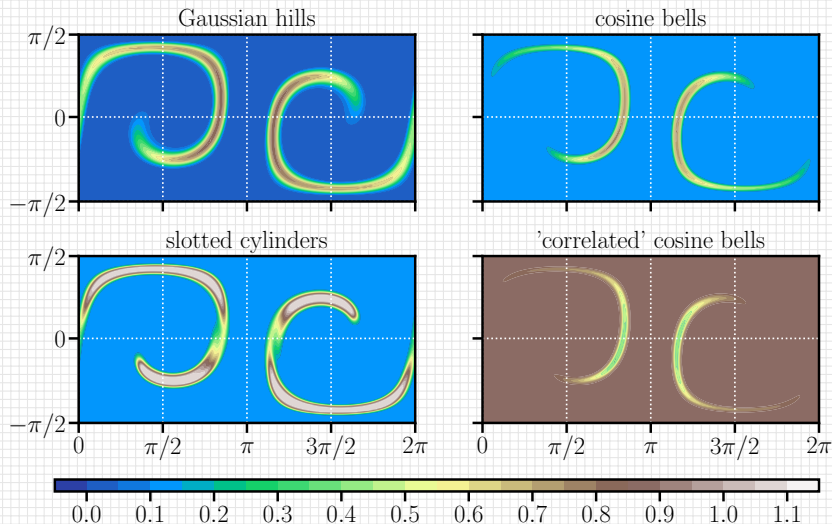
## Remarks

- Negligible difference between the fully third-order scheme and the constant-coefficient third-order variant.
- Going from second to third-order is a factor of  $\sim 3$  for the sign-preserving variant but only a factor of  $\sim 2$  for the nonoscillatory infinite-gauge variant.

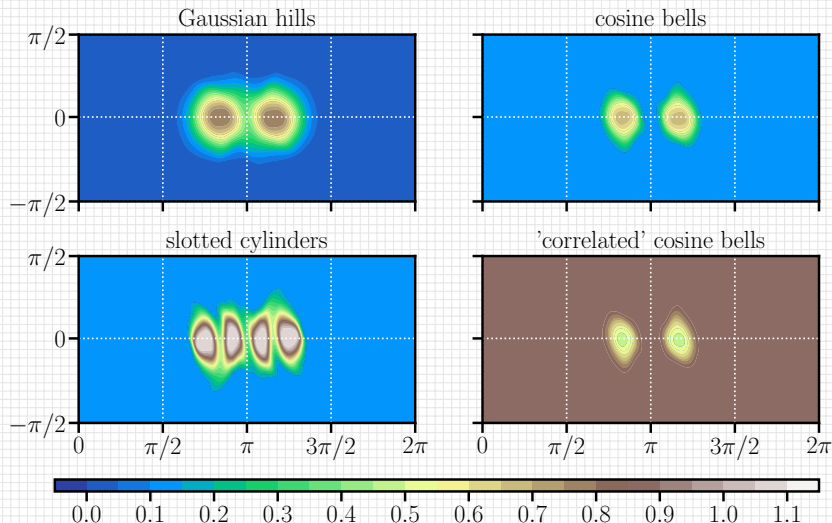
# Reversing deformational flow – Lauritzen et al. GMD 2012



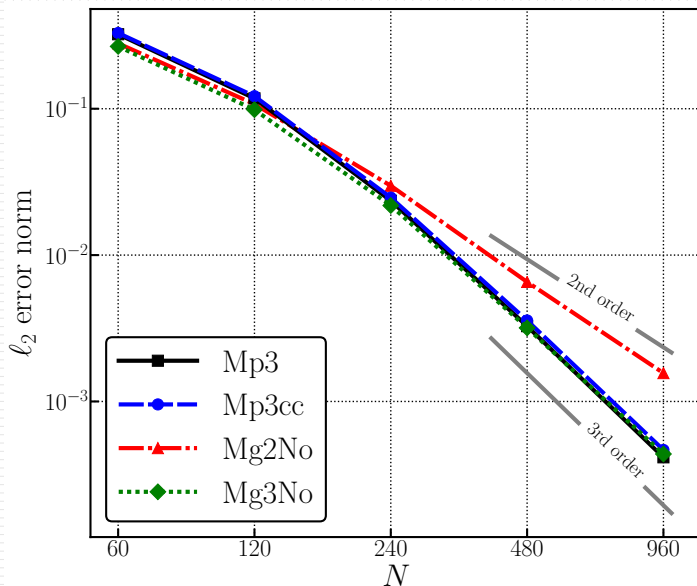
# Reversing deformational flow – Lauritzen et al. GMD 2012



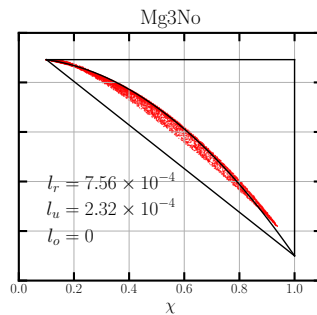
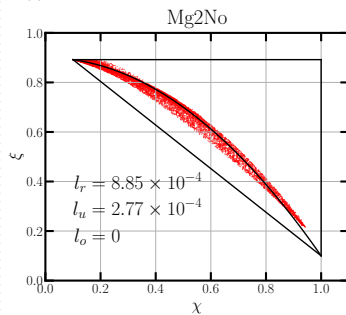
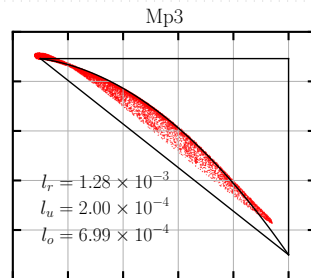
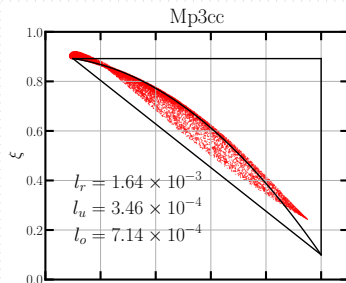
# Reversing deformational flow – Lauritzen et al. GMD 2012



# Reversing deformational flow – convergence results



# Reversing deformational flow – numerical mixing results

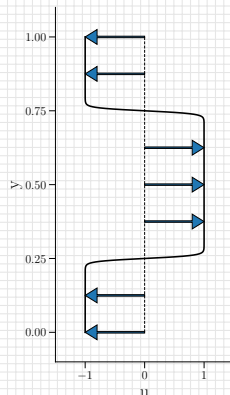


# Benefits of the fully third-order MPDATA for fluid dynamics applications

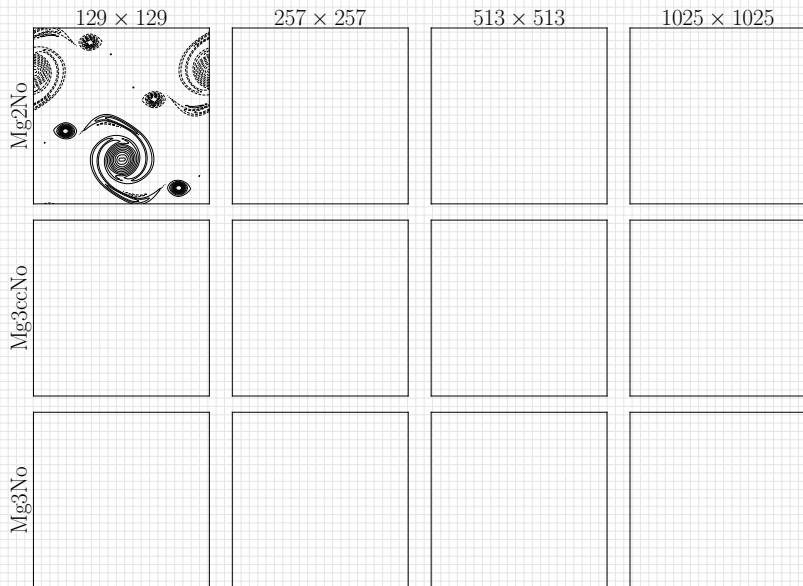


## Setup details

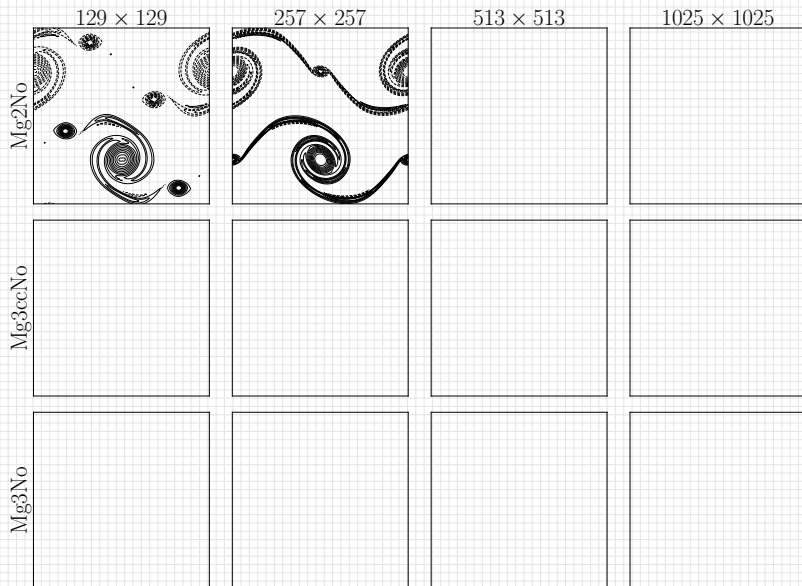
- Incompressible 2D Navier-Stokes in doubly periodic unit square.
- Only advection is integrated with third-order schemes.
- Explicit viscosity integrated to the first-order in time.
- Initial condition: shear profile  $u(y)$  as seen on the right.
- Single harmonic perturbation of  $v$  to initiate the flow.
- Reynolds number  $Re = 10000$ .



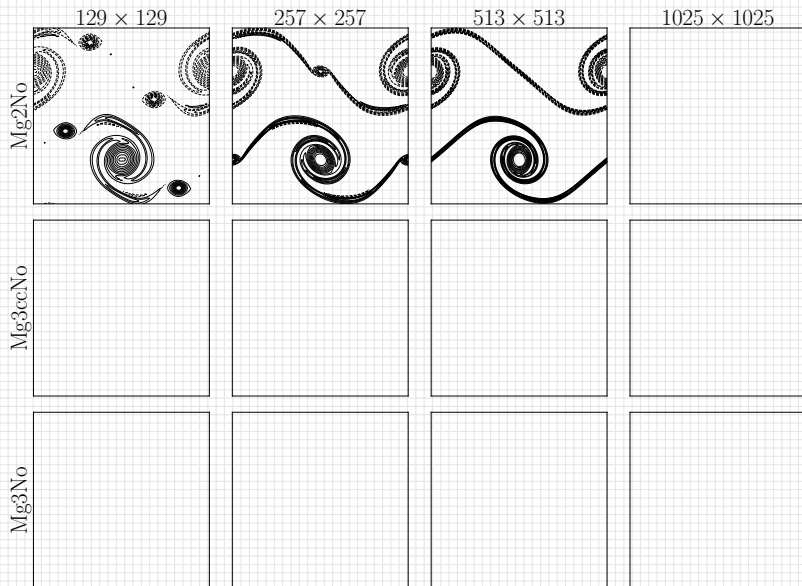
# Double shear layer rollup – vorticity contours



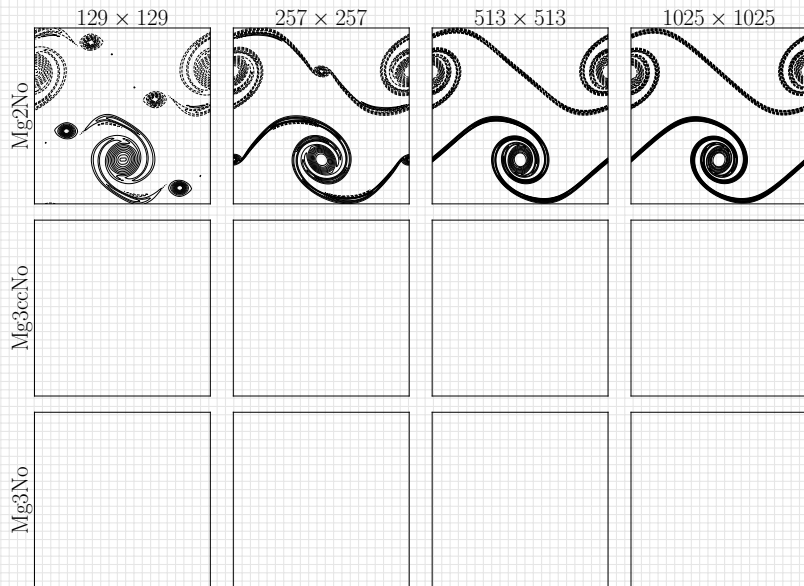
# Double shear layer rollup – vorticity contours



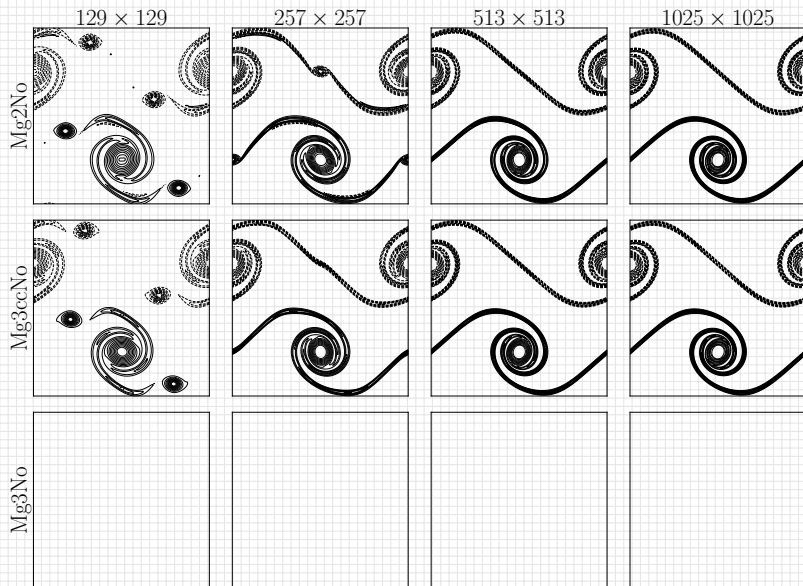
# Double shear layer rollup – vorticity contours



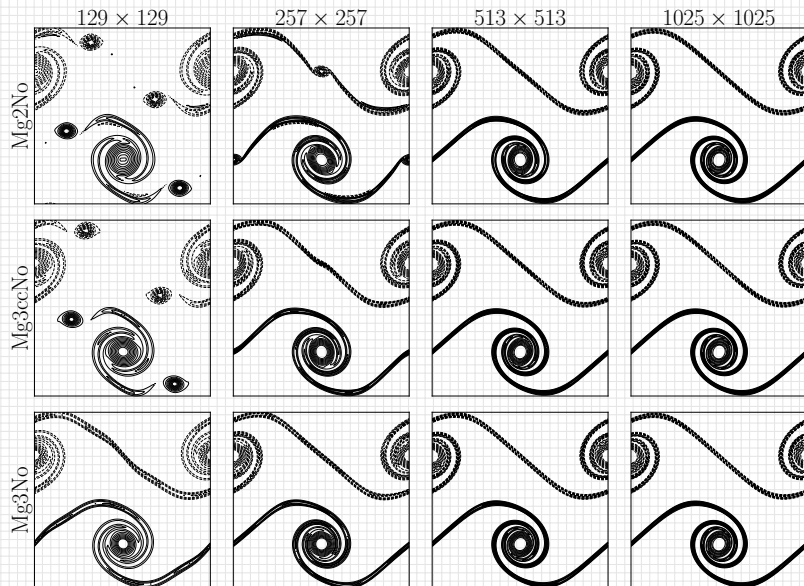
# Double shear layer rollup – vorticity contours



# Double shear layer rollup – vorticity contours



# Double shear layer rollup – vorticity contours



## $\ell_2$ error norms of $u$ after $t = 1.5$ and convergence rates

Grid	Mg2No	Order	Mg3ccNo	Order	Mg3No	Order
$129 \times 129$	$3.35 \times 10^{-1}$	—	$3.65 \times 10^{-1}$	—	$5.96 \times 10^{-2}$	—
$257 \times 257$	$1.96 \times 10^{-1}$	<b>0.77</b>	$1.09 \times 10^{-1}$	<b>1.74</b>	$4.83 \times 10^{-2}$	<b>0.30</b>
$513 \times 513$	$7.21 \times 10^{-2}$	<b>1.44</b>	$2.90 \times 10^{-2}$	<b>1.91</b>	$1.57 \times 10^{-2}$	<b>1.62</b>
$1025 \times 1025$	$2.05 \times 10^{-2}$	<b>1.82</b>	$7.06 \times 10^{-3}$	<b>2.04</b>	$4.29 \times 10^{-3}$	<b>1.87</b>

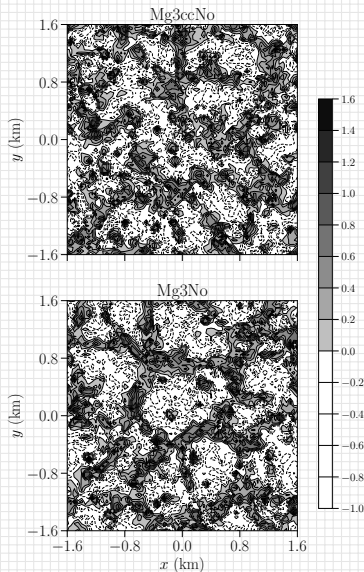
## Remarks

- Errors calculated using the Mg3No result on  $2049 \times 2049$  grid in lie of the true solution.
- $\text{Mg2No} > \text{Mg3ccNo} > \text{Mg3No}$  error ordering even for the converged solutions.
- Overall convergence rates are around 2 (not formally assured).

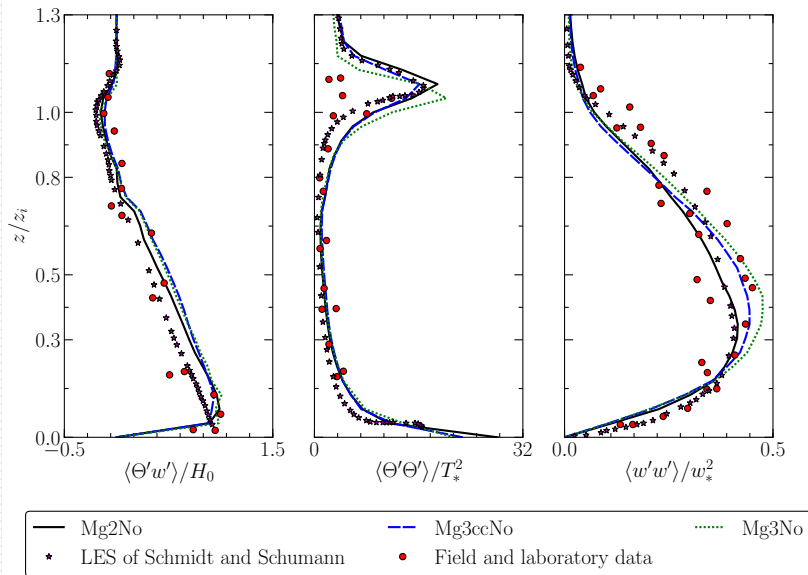


## Setup details

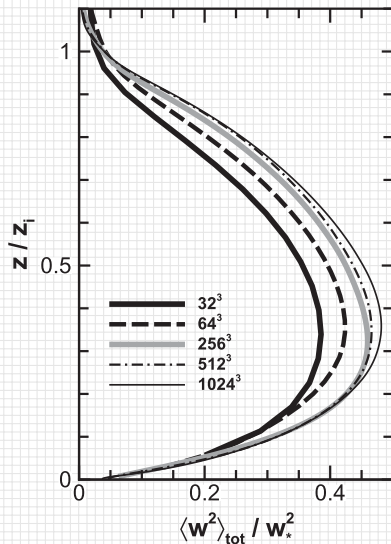
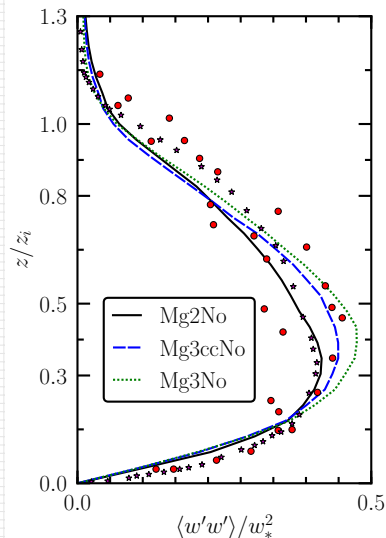
- Incompressible Boussinesq equations in 3D.
- No explicit subgrid model, ILES benchmark.
- Driven by a prescribed heat flux.
- Domain  $3.2 \text{ km} \times 3.2 \text{ km} \times 1.5 \text{ km}$ .
- Grid spacing  $50 \text{ m} \times 50 \text{ m} \times 30 \text{ m}$  ( $65 \times 65 \times 51$  points).
- Time step 8 s, simulation time  $\sim 4 \text{ h} \sim 13$  eddy turnover times.



# Convective boundary layer – profiles

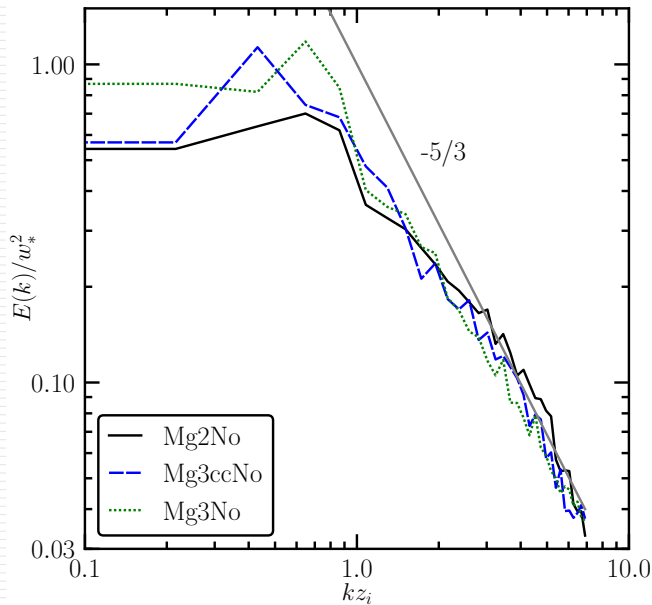


# Convective boundary layer – vertical velocity variance



Taken from Sullivan, Patton, JAS, 2011

# Convective boundary layer – spectra

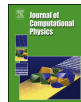




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### MPDATA: Third-order accuracy for variable flows

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#### ABSTRACT

This paper extends the multidimensional positive definite advection transport algorithm (MPDATA) to third-order accuracy for temporally and spatially varying flows. This is accomplished by identifying the leading truncation error of the standard second-order MPDATA, performing the Cauchy–Kowalevski procedure to express it in a spatial form and compensating its discrete representation—much in the same way as the standard MPDATA corrects the first-order accurate upwind scheme. The procedure of deriving the spatial form of the truncation error was automated using a computer algebra system. This enables various options in MPDATA to be included straightforwardly in the third-order scheme, thereby minimising the implementation effort in existing code bases. Following the spirit of MPDATA, the error is compensated using the upwind scheme resulting in a sign-preserving algorithm, and the entire scheme can be formulated using only two

Fully third-order MPDATA is available in *libmpdata++*

free & open source C++ library of MPDATA solvers developed in our group

*libmpdata++* repository

<https://github.com/igfuw/libmpdataxx>

Tracer transport examples from the paper

[https://github.com/igfuw/libmpdataxx/tests/mp3\\_paper\\_JCP\\_2018](https://github.com/igfuw/libmpdataxx/tests/mp3_paper_JCP_2018)

SageMath scripts used in the derivation

[https://github.com/igfuw/mpdata\\_mea](https://github.com/igfuw/mpdata_mea)

## Conclusions

- Third-order MPDATA is beneficial for tracer transport with the degree of improvement dependant on the flow and tracer field structure and measured statistics.
- Using third-order MPDATA results in improved resolution even when embedded in overall lower order accurate flow solver.
- Fully third-order MPDATA shows similar or better results than the constant-coefficient third-order variant at negligible computational expense.
- Fully third-order MPDATA can be used for ILES and the implicit subgrid model shows different scale-selectivity.

## Outlook

- Closer investigation of the implicit turbulence model.
- Extension to unstructured meshes.